

Then T is of strong type (p, p) for every $1 < p < q$ that is

$$\|Tf\|_p \leq C \|f\|_p$$

for every $f \in L^p(\mathbf{R}^n)$.

Proof. **Case** $q < \infty$. Let $f = f_1 + f_2$ where as before

$$f_1 = f\chi_{\{|f| \leq \lambda\}} \quad \text{and} \quad f_2 = f\chi_{\{|f| > \lambda\}}$$

and recall that $f_1 \in L^q$ and $f_2 \in L^1$. Subadditivity implies

$$|Tf| \leq |Tf_1| + |Tf_2|$$

for a.e. $x \in \mathbf{R}^n$. Thus

$$\begin{aligned} m(\{x \in \mathbf{R}^n : |Tf(x)| > \lambda\}) &\leq m(\{x \in \mathbf{R}^n : |Tf_1(x)| > \lambda/2\}) \\ &\quad + m(\{x \in \mathbf{R}^n : |Tf_2(x)| > \lambda/2\}) \\ &\leq \left(\frac{C_1}{\lambda/2} \|f_1\|_q \right)^q + \frac{C_2}{\lambda/2} \|f_2\|_1 \\ &\leq \frac{(2C_1)^q}{\lambda^q} \int_{\{x \in \mathbf{R}^n : |f(x)| \leq \lambda\}} |f(x)|^q \, dx \\ &\quad + \frac{2C_2}{\lambda} \int_{\{x \in \mathbf{R}^n : |f(x)| > \lambda\}} |f(x)| \, dx. \end{aligned}$$

Then by Lemma 2.9, it follows that

$$\begin{aligned} \int_{\mathbf{R}^n} |Tf|^p \, dx &= p \int_0^\infty \lambda^{p-1} m(\{x \in \mathbf{R}^n : |Tf(x)| > \lambda\}) \, d\lambda \\ &\leq (2C_1)^q p \int_0^\infty \lambda^{p-q-1} \int_{\{x \in \mathbf{R}^n : |f(x)| \leq \lambda\}} |f(x)|^q \, dx \, d\lambda \\ &\quad + 2pC_2 \int_0^\infty \lambda^{p-2} \int_{\{x \in \mathbf{R}^n : |f(x)| > \lambda\}} |f(x)| \, dx \, d\lambda. \end{aligned}$$

Further by Fubini's theorem

$$\begin{aligned} \int_0^\infty \lambda^{p-q-1} \int_{\{x \in \mathbf{R}^n : |f(x)| \leq \lambda\}} |f(x)|^q \, dx \, d\lambda &= \int_{\mathbf{R}^n} |f(x)|^q \int_{|f(x)|}^\infty \lambda^{p-q-1} \, d\lambda \, dx \\ &= \frac{1}{q-p} \int_{\mathbf{R}^n} |f(x)|^q |f(x)|^{p-q} \, dx \\ &= \frac{1}{q-p} \int_{\mathbf{R}^n} |f(x)|^p \, dx \end{aligned}$$

and

$$\begin{aligned} \int_0^\infty \lambda^{p-2} \int_{\{x \in \mathbf{R}^n : |f(x)| > \lambda\}} |f(x)| \, dx \, d\lambda &= \int_{\mathbf{R}^n} |f(x)| \int_0^{|f(x)|} \lambda^{p-2} \, d\lambda \, dx \\ &= \frac{1}{p-1} \int_{\mathbf{R}^n} |f(x)|^{p-1} |f(x)| \, dx \\ &= \frac{1}{p-1} \int_{\mathbf{R}^n} |f(x)|^p \, dx. \end{aligned}$$

Thus we arrive at

$$\|Tf\|_p^p \leq p \left(\frac{2C_2}{p-1} + \frac{(2C_1)^q}{q-p} \right) \|f\|_p^p.$$

Case $q = \infty$. Suppose that

$$\|Tg\|_\infty \leq C_2 \|g\|_\infty$$

for every $g \in L^\infty(\mathbf{R}^n)$. We again split $f \in L^p(\mathbf{R}^n)$ as $f = f_1 + f_2$ where

$$f_1 = f \chi_{\{|f| \leq \lambda/(2C_2)\}} \quad \text{and} \quad f_2 = f \chi_{\{|f| > \lambda/(2C_2)\}}$$

and by Lemma 2.20, $f_1 \in L^\infty$ and $f_2 \in L^1$. We have a.e.

$$|Tf_1(x)| \leq \|Tf_1\|_\infty \leq C_2 \|f_1\|_\infty \leq C_2 \frac{\lambda}{2C_2} = \frac{\lambda}{2}.$$

Thus

$$\begin{aligned} m(\{x \in \mathbf{R}^n : |Tf(x)| > \lambda\}) &\leq \underbrace{m(\{x \in \mathbf{R}^n : |Tf_1(x)| > \lambda/2\})}_{=0} \\ &\quad + m(\{x \in \mathbf{R}^n : |Tf_2(x)| > \lambda/2\}). \end{aligned}$$

It follows that

$$\begin{aligned} m(\{x \in \mathbf{R}^n : |Tf(x)| > \lambda\}) &\leq m(\{x \in \mathbf{R}^n : |Tf_2(x)| > \lambda/2\}) \\ &\stackrel{\text{weak } (1,1)}{\leq} \frac{C_1}{\lambda/2} \int_{\mathbf{R}^n} |f_2(x)| \, dx \\ &= \frac{2C_1}{\lambda} \int_{\{x \in \mathbf{R}^n : |f(x)| > \lambda/(2C_2)\}} |f(x)| \, dx. \end{aligned}$$

Then by using Lemma 2.9 again, we see that

$$\begin{aligned} \int_{\mathbf{R}^n} |Tf(x)|^p \, dx &= p \int_0^\infty \lambda^{p-1} m(\{x \in \mathbf{R}^n : |Tf(x)| > \lambda\}) \, d\lambda \\ &\leq 2C_1 p \int_0^\infty \lambda^{p-2} \int_{\{x \in \mathbf{R}^n : |f(x)| > \lambda/(2C_2)\}} |f(x)| \, dx \, d\lambda \\ &\stackrel{\text{Fubini}}{=} 2^p C_2^{p-1} C_1 \frac{p}{p-1} \int_{\mathbf{R}^n} |f(x)|^p \, dx. \end{aligned} \quad \square$$

Example 2.22 (Proof of the Sobolev's inequality via the maximal function). Suppose that $u \in C_0^\infty(\mathbf{R}^n)$. We immediately have

$$u(x) = - \int_0^\infty \frac{\partial}{\partial r} u(x + r\omega) dr,$$

where $\omega \in \partial B(0, 1)$. Integrating this over the whole unit sphere

$$\begin{aligned} \omega_{n-1} u(x) &= \int_{\partial B(0,1)} u(x) dS(\omega) \\ &= - \int_{\partial B(0,1)} \int_0^\infty \frac{\partial}{\partial r} u(x + r\omega) dr dS(\omega) \\ &= - \int_{\partial B(0,1)} \int_0^\infty \nabla u(x + r\omega) \cdot \omega dr dS(\omega) \\ &= - \int_0^\infty \int_{\partial B(0,1)} \nabla u(x + r\omega) \cdot \omega dS(\omega) dr \end{aligned}$$

and changing variables so that $y = x + r\omega$, $dS(y) = r^{n-1} dS(\omega)$, $\omega = (y - x)/|y - x|$, $r = |y - x|$ we get

$$\omega_{n-1} u(x) = - \int_0^\infty \int_{\partial B(0,r)} \nabla u(y) \cdot \frac{y - x}{|y - x|^n} dS(y) dr$$

so that

$$u(x) = - \frac{1}{\omega_{n-1}} \int_{\mathbf{R}^n} \frac{\nabla u(y) \cdot (x - y)}{|x - y|^n} dy.$$

Further

$$|u(x)| \leq \frac{1}{\omega_{n-1}} \int_{\mathbf{R}^n} \frac{|\nabla u(y)|}{|x - y|^{n-1}} dy$$

which is so called Riesz potential. We split this into a bad part and a good part as $\int_{\mathbf{R}^n} = \int_{B(x,r)} + \int_{\mathbf{R}^n \setminus B(x,r)}$. By estimating the bad part over the sets $B(x, 2^{-i}r) \setminus B(x, 2^{-i-1}r)$ as

$$\begin{aligned} \int_{B(x,r)} \frac{|\nabla u(y)|}{|x - y|^{n-1}} dy &= \sum_{i=0}^{\infty} \int_{B(x, 2^{-i}r) \setminus B(x, 2^{-i-1}r)} \frac{|\nabla u(y)|}{|x - y|^{n-1}} dy \\ &\leq \sum_{i=0}^{\infty} \int_{B(x, 2^{-i}r) \setminus B(x, 2^{-i-1}r)} \frac{|\nabla u(y)|}{(2^{-i-1}r)^{n-1}} dy \\ &\leq \sum_{i=0}^{\infty} \frac{2^{-i}r}{2^{-i}r} \int_{B(x, 2^{-i}r)} 2^{n-1} \frac{|\nabla u(y)|}{(2^{-i}r)^{n-1}} dy \\ &\leq C \sum_{i=0}^{\infty} 2^{n-1} 2^{-i} r \int_{B(x, 2^{-i}r)} |\nabla u(y)| dy \\ &\leq C 2^{n-1} r M |\nabla u|(x) \sum_{i=0}^{\infty} 2^{-i} \end{aligned}$$

we get

$$\int_{B(x,r)} \frac{|\nabla u(y)|}{|x-y|^{n-1}} dy \leq CrM |\nabla u|(x). \quad (2.23)$$

On the other hand, for the good part we use Hölder's inequality with the powers p and $p/(p-1)$, where $p < n$, as

$$\begin{aligned} & \int_{\mathbf{R}^n \setminus B(x,r)} \frac{|\nabla u(y)|}{|x-y|^{n-1}} dy \\ & \leq \left(\int_{\mathbf{R}^n \setminus B(x,r)} |\nabla u(y)|^p dy \right)^{1/p} \left(\int_{\mathbf{R}^n \setminus B(x,r)} |x-y|^{(1-n)p/(p-1)} dy \right)^{(p-1)/p}. \end{aligned}$$

Then we calculate

$$\begin{aligned} & \left(\int_{\mathbf{R}^n \setminus B(x,r)} |x-y|^{(1-n)p/(p-1)} dy \right)^{(p-1)/p} \\ & = \left(\int_r^\infty \omega_{n-1} \rho^{n-1} \rho^{(1-n)p/(p-1)} d\rho \right)^{(p-1)/p} \\ & = \left(\omega_{n-1} \int_r^\infty \rho^{(1-n)/(p-1)} d\rho \right)^{(p-1)/p} = \left(\omega_{n-1} \int_r^\infty \rho^{-1+(p-n)/(p-1)} d\rho \right)^{(p-1)/p}. \end{aligned}$$

Combining the previous calculations, we get

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$$\int_{\mathbf{R}^n \setminus B(x,r)} \frac{|\nabla u(y)|}{|x-y|^{n-1}} dy \leq C \|\nabla u\|_p r^{1-\frac{n}{p}}, \quad (2.24)$$

with $p < n$. Choosing $r = (\|\nabla u\|_p / (M |\nabla u|(x)))^{p/n}$ as well as combining the estimates (2.23) and (2.24), we get

$$\begin{aligned} |u(x)| & \leq C \int_{\mathbf{R}^n} \frac{|\nabla u(y)|}{|x-y|^{n-1}} dy \\ & \leq C \|\nabla u\|_p^{p/n} M |\nabla u|(x)^{(n-p)/n}. \end{aligned}$$

Then we take the power¹ $np/(n-p)$ on both sides and end up with

$$|u(x)|^{np/(n-p)} \leq C \|\nabla u\|_p^{p^2/(n-p)} M |\nabla u|(x)^p.$$

By recalling Hardy-Littlewood II, we obtain

$$\begin{aligned} \int_{\mathbf{R}^n} |u(x)|^{np/(n-p)} dx & \leq C \|\nabla u\|_p^{p^2/(n-p)} \int_{\mathbf{R}^n} M |\nabla u|(x)^p dx \\ & \leq C \|\nabla u\|_p^{p^2/(n-p)} \|\nabla u\|_p^p \leq C \|\nabla u\|_p^{np/(n-p)}. \end{aligned}$$

This is so called Sobolev's inequality

$$\left(\int_{\mathbf{R}^n} |u(x)|^{p^*} dx \right)^{1/p^*} \leq C \left(\int_{\mathbf{R}^n} |\nabla u(x)|^p dx \right)^{1/p},$$

which holds for every $u \in C_0^\infty(\mathbf{R}^n)$ and $p < n$.

¹This is sometimes denoted by $p^* = np/(n-p)$ and called a Sobolev conjugate. It satisfies $1/p + 1/p^* = 1/n$.