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#### Abstract

Discrete maximum principles are proved for finite element discretizations of nonlinear elliptic interface problems with jumps of the normal derivatives. The geometric conditions in the case of simplicial meshes are suitable acuteness or nonobtuseness properties.


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## 1 Introduction

The maximum principle forms an important qualitative property of second order elliptic boundary value problems [10, 23, 27]. Consequently, the discrete analogues of the maximum principle (so-called discrete maximum principles, DMPs) have drawn much attention. Various DMPs have been formulated and proved including the case of finite difference, finite volume and finite element approximations, and corresponding geometric conditions on the computational meshes have been given, see, e.g., $[2,4,6,7,11,19,28]$ for linear and $[14,15,20]$ for nonlinear problems with standard (i.e., Dirichlet, and in $[14,15]$ mixed) boundary conditions.

In this paper we address interface problems, which arise in various branches of material science, biochemistry, multiphase flow etc., often when two distinct materials are involved with different conductivities or densities. Another (for our work, motivating) example is from localized reaction-diffusion problems $[12,13]$, see at the end of this paper. Many special numerical methods have been designed for interface problems, see e.g. [12, 25, 26, 24], but maximum principles have received less attention than for the case of standard boundary value problems. A continuous minimum principle for a related problem is given in [9]. The discrete maximum principle for suitable finite difference discretizations of linear interface problems has been proved in [25].

Our goal is to prove discrete maximum principles for nonlinear elliptic interface problems when finite element discretization is involved. The present paper is the extension of our paper [14] to a class of such problems, and relies on a similar technique using weak formulation and positivity conditions that ensure well-posedness. We consider matching conditions for the solution itself on the interface, i.e., the jump is allowed for the normal derivatives. Problems with jump of the solution or without well-posedness may be the subject of further research.

The paper is organized as follows. The formulation of the problem, together with the derivation of a continuous maximum principle, and the description of finite element discretization are given in Section 2. Discrete maximum principles are proved and examples are given in Section 3.

## 2 Nonlinear elliptic interface problems: basic properties and discretization

### 2.1 Formulation of the problem

We investigate nonlinear interface problems of the following form:

$$
\left\{\begin{align*}
-\operatorname{div}(b(x, \nabla u) \nabla u)+q(x, u) & =f(x) & & \text { in } \Omega \backslash \Gamma  \tag{1}\\
{[u]_{\Gamma} } & =0 & & \text { on } \Gamma \\
{\left[b(x, \nabla u) \frac{\partial u}{\partial \nu}\right]_{\Gamma}+s(x, u) } & =\gamma(x) & & \text { on } \Gamma \\
u & =g(x) & & \text { on } \partial \Omega,
\end{align*}\right.
$$

where $\partial \Omega$ denotes the boundary of the domain $\Omega$ and the interface $\Gamma$ is a surface lying in $\Omega$, further, $[u]_{\Gamma}$ and $\left[b(x, \nabla u) \frac{\partial u}{\partial \nu}\right]_{\Gamma}$ denote the jump (i.e., the difference of the limits from the two sides of the interface $\Gamma$ ) of $u$ and $b(x, \nabla u) \frac{\partial u}{\partial \nu}$, respectively. We impose the following

## Assumptions 2.1:

(A1) $\Omega$ is a bounded open domain in $\mathbf{R}^{d}$, the interface $\Gamma \subset \Omega$ and the boundary $\partial \Omega$ are piecewise smooth and Lipschitz continuous $(d-1)$ dimensional surfaces.
(A2) The scalar functions $b: \Omega \times \mathbf{R}^{d} \rightarrow \mathbf{R}, \quad q: \Omega \times \mathbf{R} \rightarrow \mathbf{R}$ and $s:$ $\Gamma \times \mathbf{R} \rightarrow \mathbf{R}$ are measurable and bounded w.r.t. their first variable $x \in \Omega$ (resp. $x \in \Gamma$ ) and continuously differentiable w.r.t. their second variable $\eta \in \mathbf{R}^{d}$ (resp. $\xi \in \mathbf{R}$ ). Further, $f \in L^{2}(\Omega), \gamma \in L^{2}(\Gamma)$ and $g \in H^{1}(\Omega)$.
(A3) The function $b$ satisfies

$$
\begin{equation*}
0<\mu_{0} \leq b(x, \eta) \leq \mu_{1} \tag{2}
\end{equation*}
$$

with positive constants $\mu_{0}$ and $\mu_{1}$ independent of $(x, \eta)$, further, the diadic product matrix $\eta \cdot \frac{\partial b(x, \eta)}{\partial \eta}$ is symmetric positive semidefinite and bounded in matrix norm by some positive constant $\mu_{2}$ independent of $(x, \eta)$.
(A4) Let $2 \leq p_{1}$ if $d=2$, or $2 \leq p_{1} \leq \frac{2 d}{d-2}$ if $d>2$, further, let $2 \leq p_{2}$ if $d=2$, or $2 \leq p_{2} \leq \frac{2 d-2}{d-2}$ if $d>2$. There exist functions $\alpha_{1} \in L^{d / 2}(\Omega)$, $\alpha_{2} \in L^{d-1}(\Gamma)$ and a constant $\beta \geq 0$ such that for any $x \in \Omega$ (or $x \in \Gamma$, resp.) and $\xi \in \mathbf{R}$

$$
\begin{equation*}
0 \leq \frac{\partial q(x, \xi)}{\partial \xi} \leq \alpha_{1}(x)+\beta|\xi|^{p_{1}-2}, \quad 0 \leq \frac{\partial s(x, \xi)}{\partial \xi} \leq \alpha_{2}(x)+\beta|\xi|^{p_{2}-2} \tag{3}
\end{equation*}
$$

Remark 1. The role of assumption (A3) is to ensure that the Jacobian matrices $J(x, \eta):=\frac{\partial}{\partial \eta}(b(x, \eta) \eta)$ are symmetric and satisfy the uniform ellipticity property $\mu_{0}|\zeta|^{2} \leq J(x, \eta) \zeta \cdot \zeta \leq \mu_{3}|\zeta|^{2}, \quad \zeta \in \mathbf{R}^{d}$ (with $\mu_{3}=$ $\mu_{1}+\mu_{2}$ ), which will be required for well-posedness. For instance, assumption (A3) holds for coefficients of the form $b(x, \eta)=a(x,|\eta|)$ (see [8, 21] for such nonlinearities), where the $C^{1}$ function $a: \Omega \times \mathbf{R}^{+} \rightarrow \mathbf{R}$ satisfies
$0<\mu_{0} \leq a(x, r) \leq \frac{\partial}{\partial r}(a(x, r) r) \leq \mu_{3}(r>0)$. More specially, one may have $b(x, \eta)=a(x)$ (i.e., linear principal part) with a measurable function $a$ satisfying $0<\mu_{0} \leq a(x) \leq \mu_{3}$.

### 2.2 Weak solutions

Theorem 1. Under Assumptions 2.1, problem (1) has a unique weak solution $u^{*} \in H^{1}(\Omega)$ defined as follows:

$$
\begin{gather*}
\int_{\Omega}\left(b\left(x, \nabla u^{*}\right) \nabla u^{*} \cdot \nabla v+q\left(x, u^{*}\right) v\right) d x+\int_{\Gamma} s\left(x, u^{*}\right) v d \sigma= \\
=\int_{\Omega} f v d x+\int_{\Gamma} \gamma v d \sigma \quad \forall v \in H_{0}^{1}(\Omega)  \tag{4}\\
\text { and } \quad u^{*}=g \quad \text { on } \partial \Omega . \tag{5}
\end{gather*}
$$

Proof. We first prove the theorem for homogeneous boundary condition, i.e. when $g=0$. In this case the weak solution $u^{*}$ can be obtained using monotone operators, in a similar way as in [8, Chap. 6], therefore we only indicate the main steps of the proof. First, we define

$$
\begin{gather*}
\langle F(u), v\rangle=\int_{\Omega}(b(x, \nabla u) \nabla u \cdot \nabla v+q(x, u) v-f v) d x+ \\
\quad+\int_{\Gamma}(s(x, u) v-\gamma v) d \sigma \quad\left(v \in H_{0}^{1}(\Omega)\right) \tag{6}
\end{gather*}
$$

where the growth conditions in (A1)-(A4) ensure that the arising integrals are finite. Let $J(x, \eta):=\frac{\partial}{\partial \eta}(b(x, \eta) \eta)$ as in Remark 1. Then, from (A3)(A4), we obtain that the Gateaux derivative $F^{\prime}(u)$ exists, is self-adjoint for all $u$ and satisfies

$$
\begin{align*}
\left\langle F^{\prime}(u) v, v\right\rangle & =\int_{\Omega}\left(J(x, \nabla u) \nabla v \cdot \nabla v+q_{u}^{\prime}(x, u) v^{2}\right) d x+ \\
& +\int_{\Gamma} s_{u}^{\prime}(x, u) v^{2} d \sigma \geq \mu_{0} \int_{\Omega}|\nabla v|^{2} d x \tag{7}
\end{align*}
$$

(for all $u, v \in H_{0}^{1}(\Omega)$ ), where $q_{u}^{\prime}$, $s_{u}^{\prime}$ denote derivatives w.r.t. $u$. Using the standard Sobolev norm defined via

$$
\|v\|_{1}^{2}=\int_{\Omega}|\nabla v|^{2} d x \quad\left(v \in H_{0}^{1}(\Omega)\right)
$$

the uniform ellipticity (7) implies that the operator equation $F(u)=0$ has a unique solution $u^{*} \in H_{0}^{1}(\Omega)$. Here $F\left(u^{*}\right)=0$ is equivalent to (4), i.e., $u^{*}$ is the weak solution.

For non-homogeneous boundary conditions the problem can be reduced to the homogeneous case using a usual translation. Let $g \in H^{1}(\Omega)$ be arbitrary and let us require (5) on the boundary. Then we look for $u^{*}$ in the form $u^{*}=u+g$, in which case $u=0$ on $\partial \Omega$. Substituting this sum into (4), we observe that $u$ must satisfy the same problem with homogeneous boundary conditions and with coefficients
$\hat{b}(x, \eta)=b(x, \eta+\nabla g(x)), \quad \hat{q}(x, \xi)=q(x, \xi+g(x)), \quad \hat{s}(x, \xi)=s(x, \xi+g(x))$.
Here $g(x)$ is independent of $\xi, \eta$, hence these coefficients remain $C^{1}$ in their second variable and satisfy the same growth conditions as $b, q, s$. This implies existence and uniqueness for $u$, and then the same for $u^{*}$ owing to the relation $u^{*}=u+g$.

The above notion of weak solution is justified by showing that any classical (or strong) solution is also a weak solution. To define the classical solution, we assume in addition that the interface $\Gamma$ is a closed surface, or more generally, it is any compact subset of an (also piecewise smooth and Lipschitz continuous) closed surface $\hat{\Gamma} \subset \Omega$ as illustrated in Figure 2.2. Let us denote by $\Omega_{0}$ the interior of the surface $\hat{\Gamma}$.


Figure 1: Interface in a domain

Definition 1. We call $u: \bar{\Omega} \rightarrow \mathbf{R}$ a classical solution of problem (1) if $u \in C^{2}(\Omega \backslash \Gamma),\left.\quad u\right|_{\bar{\Omega}_{0}} \in C^{1}\left(\bar{\Omega}_{0}\right),\left.\quad u\right|_{{\bar{\Omega} \backslash \Omega_{0}} \in C^{1}\left({\bar{\Omega} \backslash \Omega_{0}}\right) \text { and } u \text { satisfies (1) }}$ pointwise.

Proposition 1. A classical solution of problem (1) is also a weak solution.
Proof. Let $x \in \hat{\Gamma}$, let $\nu$ denote the normal unit vector pointing out of $\Omega_{0}$, and let $\hat{\nu}:=-\nu$ (the normal unit vector pointing out of $\Omega \backslash \Omega_{0}$ ). The jump of $b(x, \nabla u) \frac{\partial u}{\partial \nu}$ at $x$ is the difference of the limits of $b(\cdot, \nabla u) \frac{\partial u}{\partial \nu}$ at $x$ from $\Omega_{0}$ and from $\Omega \backslash \Omega_{0}$. Using the definition $\frac{\partial u}{\partial \nu}(x):=\lim _{t \rightarrow 0^{+}} \frac{1}{t}(u(x)-u(x-t \nu))$, we thus have

$$
\left[b(x, \nabla u) \frac{\partial u}{\partial \nu}\right]_{x \in \Gamma}:=b\left(x, \nabla\left(\left.u\right|_{\bar{\Omega}_{0}}\right)(x)\right) \lim _{t \rightarrow 0^{+}} \frac{1}{t}(u(x)-u(x-t \nu))-
$$

$$
\begin{gather*}
-b\left(x, \nabla\left(\left.u\right|_{\overline{\Omega \backslash \Omega_{0}}}\right)(x)\right) \lim _{t \rightarrow 0^{-}} \frac{1}{t}(u(x)-u(x-t \nu))= \\
=b\left(x, \nabla\left(\left.u\right|_{\bar{\Omega}_{0}}\right)(x)\right) \lim _{t \rightarrow 0^{+}} \frac{1}{t}(u(x)-u(x-t \nu))+ \\
+b\left(x, \nabla\left(\left.u\right|_{\overline{\Omega \backslash \Omega_{0}}}\right)(x)\right) \lim _{s \rightarrow 0^{+}} \frac{1}{s}(u(x)-u(x-s \hat{\nu}))= \\
=\left(b\left(x, \nabla\left(\left.u\right|_{\bar{\Omega}_{0}}\right)(x)\right) \frac{\partial u}{\partial \nu}+b\left(x, \nabla\left(\left.u\right|_{\overline{\Omega \backslash \Omega_{0}}}\right)(x)\right) \frac{\partial u}{\partial \hat{\nu}}\right)_{x \in \Gamma} . \tag{8}
\end{gather*}
$$

Now let $u$ be a classical solution. The assumptions imply that $u \in H^{1}(\Omega)$, and (5) holds trivially. For any $v \in H_{0}^{1}(\Omega)$, Green's formula for equation (1) on $\Omega_{0}$ and $\Omega \backslash \Omega_{0}$, respectively, yields

$$
\int_{\Omega_{0}} f v d x=\int_{\Omega_{0}}(b(x, \nabla u) \nabla u \cdot \nabla v+q(x, u) v) d x-\int_{\hat{\Gamma}} b\left(x, \nabla\left(\left.u\right|_{\bar{\Omega}_{0}}\right)\right) \frac{\partial u}{\partial \nu} v d \sigma
$$

and
$\int_{\Omega \backslash \Omega_{0}} f v d x=\int_{\Omega \backslash \Omega_{0}}(b(x, \nabla u) \nabla u \cdot \nabla v+q(x, u) v) d x-\int_{\hat{\Gamma}} b\left(x, \nabla\left(\left.u\right|_{\bar{\Omega} \Omega_{0}}\right)\right) \frac{\partial u}{\partial \hat{\nu}} v d \sigma$.
Summing up, the integrand on $\hat{\Gamma}$ becomes the jump on $\Gamma$ (using (8)) and zero on $\hat{\Gamma} \backslash \Gamma$ (since $\nabla u$ is continuous there and $\hat{\nu}=-\nu$ ). In virtue of the jump condition in (1), we altogether obtain

$$
\begin{gathered}
\int_{\Omega} f v d x=\int_{\Omega}(b(x, \nabla u) \nabla u \cdot \nabla v+q(x, u) v) d x-\int_{\Gamma}\left[b(x, \nabla u) \frac{\partial u}{\partial \nu}\right]_{\Gamma} v d \sigma \\
=\int_{\Omega}(b(x, \nabla u) \nabla u \cdot \nabla v+q(x, u) v) d x+\int_{\Gamma}((s(x, u)-\gamma) v d \sigma
\end{gathered}
$$

### 2.3 Continuous maximum principles

We formulate and prove two continuous maximum principles for our PDE problem (1). These statements provide the properties whose discrete analogues can be expected for suitable FEM solutions.
Theorem 2. Let Assumptions 2.1 hold and

$$
\begin{equation*}
f(x)-q(x, 0) \leq 0, x \in \Omega, \quad \text { and } \quad \gamma(x)-s(x, 0) \leq 0, x \in \Gamma \tag{9}
\end{equation*}
$$

If the weak solution $u$ of problem (1) belongs to $C^{1}(\Omega \backslash \Gamma) \cap C(\bar{\Omega})$, then

$$
\begin{equation*}
\max _{\bar{\Omega}} u \leq \max \left\{0, \max _{\partial \Omega} g\right\} . \tag{10}
\end{equation*}
$$

In particular, if $g \geq 0$, then $\max _{\bar{\Omega}} u=\max _{\partial \Omega} g$, and if $g \leq 0$, then we have the nonpositivity property $\max _{\bar{\Omega}} u \leq 0$.

In general, if $u \in H^{1}(\Omega)$ only (without the above regularity assumption) and $g$ is a.e. bounded on $\partial \Omega$, then the same statements hold if $\max u$ and $\max g$ are replaced by ess sup $u$ and ess sup $g$, respectively.

Proof. We only prove the regular case, the general case is similar (if $\max u$ and $\max g$ are replaced by ess $\sup u$ and $\operatorname{ess} \sup g$, respectively). Let

$$
r(x, \xi):=\left\{\begin{array}{ll}
\frac{q(x, \xi)-q(x, 0)}{\xi}, & \text { if } \xi \neq 0,  \tag{11}\\
\frac{\partial q}{\partial \xi}(x, 0), & \text { if } \xi=0,
\end{array} \quad z(x, \xi):= \begin{cases}\frac{s(x, \xi)-s(x, 0)}{\xi}, & \text { if } \xi \neq 0 \\
\frac{\partial s}{\partial \xi}(x, 0), & \text { if } \xi=0 .\end{cases}\right.
$$

Here, using (A2), the functions $r$ and $z$ are continuous in $\xi$. Further, in view of (A4), we have

$$
\begin{equation*}
r(x, \xi) \geq 0, \quad z(x, \xi) \geq 0 \tag{12}
\end{equation*}
$$

We define

$$
\begin{gather*}
\tilde{a}(x):=b(x, \nabla u(x)) \quad(x \in \Omega \backslash \Gamma), \quad \tilde{h}(x):=r(x, u(x)) \quad(x \in \Omega),  \tag{13}\\
\tilde{k}(x):=z(x, u(x)) \quad(x \in \Gamma) .
\end{gather*}
$$

Using also the notations

$$
\begin{equation*}
\hat{f}(x):=f(x)-q(x, 0) \quad \text { and } \quad \hat{\gamma}(x):=\gamma(x)-s(x, 0), \tag{14}
\end{equation*}
$$

the weak formulation of problem (1) is rewritten as

$$
\begin{equation*}
\int_{\Omega}(\tilde{a} \nabla u \cdot \nabla v+\tilde{h} u v) d x+\int_{\Gamma} \tilde{k} u v d \sigma=\int_{\Omega} \hat{f} v d x+\int_{\Gamma} \hat{\gamma} v d \sigma \quad \forall v \in H_{0}^{1}(\Omega) . \tag{15}
\end{equation*}
$$

Let $M:=\max \left\{0, \max _{\partial \Omega} g\right\}$ and we introduce the piecewise $C^{1}$ function $v:=$ $\max \{u-M, 0\}$. Then we have $v \geq 0$ and $v_{\mid \partial \Omega}=0$, further, $u(x)=v(x)+M$ for any $x \in \bar{\Omega}$ unless $v(x)=0$. Hence, for this $v$ the left-hand side of (15) satisfies

$$
\begin{gathered}
\int_{\Omega}(\tilde{a} \nabla u \cdot \nabla v+\tilde{h} u v) d x+\int_{\Gamma} \tilde{k} u v d \sigma= \\
=\int_{\Omega}\left(\tilde{a}|\nabla v|^{2}+\tilde{h}(v+M) v\right) d x+\int_{\Gamma} \tilde{k}(v+M) v d \sigma \geq 0
\end{gathered}
$$

since the functions $\tilde{a}, \tilde{h}, \tilde{k}, v$ and the constant $M$ are nonnegative. On the other hand, the assumptions $\hat{f} \leq 0, \hat{\gamma} \leq 0$ imply that for this $v$ the righthand side of (15) satisfies

$$
\int_{\Omega} \hat{f} v d x+\int_{\Gamma} \hat{\gamma} v d \sigma \leq 0
$$

which together imply the relation

$$
\int_{\Omega}\left(\tilde{a}|\nabla v|^{2}+\tilde{h}(v+M) v\right) d x+\int_{\Gamma} \tilde{k}(v+M) v d \sigma=0 .
$$

By assumption (A3), here $\tilde{a}$ has a positive minimum, hence $|\nabla v|=0$, i.e., $v$ is constant. We have seen that $v_{\mid \partial \Omega}=0$, hence we obtain that $v \equiv 0$, which just means that (10) holds.

The following special case provides equality of maxima on $\partial \Omega$ without assuming $g \geq 0$ :

Theorem 3. Let $q \equiv 0$ and $s \equiv 0$ in problem (1). Let us impose the assumptions of Theorem 2, which now means that (A1)-(A3) are satisfied, $u \in C^{1}(\Omega \backslash \Gamma) \cap C(\bar{\Omega})$, and (9) takes the form

$$
\begin{equation*}
f(x) \leq 0, x \in \Omega \quad \text { and } \quad \gamma(x) \leq 0, x \in \Gamma . \tag{16}
\end{equation*}
$$

Then

$$
\begin{equation*}
\max _{\bar{\Omega}} u=\max _{\partial \Omega} g . \tag{17}
\end{equation*}
$$

(If $u \in H^{1}(\Omega)$ only and $g$ is a.e. bounded on $\partial \Omega$, then ess $\sup u=\operatorname{ess} \sup g$ on $\partial \Omega$.)

Proof. We only prove the regular case again. If $\max _{\partial \Omega} g \geq 0$ then (10) implies (17). Let $\max _{\partial \Omega} g<0$, say, $\max _{\partial \Omega} g=-K$ with some $K>0$. Then the function $w:=u+K$ satisfies the same mixed problem with right-hand sides $f, \gamma$ and $g+K$, respectively, hence Theorem 2 is valid for this problem as well, and (10) for $w$ yields $\max _{\bar{\Omega}} w \leq \max \left\{0, \max _{\partial \Omega}(g+K)\right\}=0$. Then $\max _{\bar{\Omega}} u \leq-K=\max _{\partial \Omega} g$.

Remark 2. Analogously to Theorems 2 and 3, corresponding minimum principles and nonnegativity property hold if the sign conditions in (9) and (16) are reversed.

### 2.4 Finite element discretization

Our basic assumption in the sequel that $\Omega$ is a polytopic domain and the interface $\Gamma$ is also polytopic. (We note that if $\partial \Omega$ or $\Gamma$ are curved then the convergence of the discrete solution to the exact one is a much more difficult problem, out of the scope of this paper. Even for the simpler case of Dirichlet problems in 3D without interface, such an analysis has been given only recently in [16].)

We introduce a finite element discretization of our problem with simplicial elements and continuous piecewise linear basis functions. Let $\mathcal{T}_{h}$ be a a conforming triangulation of $\bar{\Omega}$ into tetrahedra, whose nodes are $B_{1}, \ldots, B_{\bar{n}}$. Denote by $\phi_{1}, \ldots, \phi_{\bar{n}}$ the piecewise linear continuous basis functions defined in a standard way, i.e., $\phi_{i}\left(B_{j}\right)=\delta_{i j}$ for $i, j=1, \ldots, \bar{n}$, where $\delta_{i j}$ is the Kronecker symbol. Let $V_{h}$ denote the finite element subspace spanned by the above basis functions:

$$
V_{h}=\operatorname{span}\left\{\phi_{1}, \ldots, \phi_{\bar{n}}\right\} \subset H^{1}(\Omega) .
$$

Let $n<\bar{n}$ be such that

$$
\begin{equation*}
B_{1}, \ldots, B_{n} \tag{18}
\end{equation*}
$$

are the nodes that lie in $\Omega$ and let

$$
\begin{equation*}
B_{n+1}, \ldots, B_{\bar{n}} \tag{19}
\end{equation*}
$$

be the nodes that lie on $\partial \Omega$. Then the basis functions $\phi_{1}, \ldots, \phi_{n}$ satisfy homogeneous boundary condition on $\partial \Omega$, i.e., $\phi_{i} \in H_{0}^{1}(\Omega)$. We define

$$
V_{h}^{0}=\operatorname{span}\left\{\phi_{1}, \ldots, \phi_{n}\right\} \subset H_{0}^{1}(\Omega) .
$$

Further, let

$$
\begin{equation*}
g_{h}=\sum_{j=n+1}^{\bar{n}} g_{j} \phi_{j} \in V_{h} \tag{20}
\end{equation*}
$$

(with $g_{j} \in \mathbf{R}$ ) be the piecewise linear approximation of the function $g$ on $\partial \Omega$ (and on the neighbouring elements). To find the FEM solution of (4)-(5) in $V_{h}$, we solve the following problem: find $u_{h} \in V_{h}$ such that

$$
\begin{gather*}
\int_{\Omega}\left(b\left(x, \nabla u_{h}\right) \nabla u_{h} \cdot \nabla v_{h}+q\left(x, u_{h}\right) v_{h}\right) d x+\int_{\Gamma} s\left(x, u_{h}\right) v_{h} d \sigma= \\
=\int_{\Omega} f v_{h} d x+\int_{\Gamma} \gamma v_{h} d \sigma \quad \forall v_{h} \in V_{h}^{0}  \tag{21}\\
\text { and } \quad u_{h}=g_{h} \quad \text { on } \partial \Omega .
\end{gather*}
$$

Theorem 4. Under Assumptions 2.1, problem (21) has a unique solution $u_{h} \in V_{h}$, and $\left\|u^{*}-u_{h}\right\|_{1} \rightarrow 0$ as $h \rightarrow 0$.

Proof. The proof of Theorem 1 can be repeated to obtain $u_{h}$, just replacing $H^{1}(\Omega)$ by $V_{h}$. The convergence of $u_{h}$ to $u^{*}$ in $H^{1}$-norm follows in the standard way from the ellipticity of the equation and the fact that the finite-dimesional subspaces $V_{h}$ satisfy the condition $\lim _{h \rightarrow 0} \operatorname{dist}\left(u, V_{h}\right)=0$ for all $u \in H^{1}(\Omega)$, where $\operatorname{dist}\left(u, V_{h}\right)=\inf _{v_{h} \in V_{h}}\left\|u-v_{h}\right\|_{1} \quad$ (see [5]).

Let us now formulate the nonlinear algebraic system corresponding to (21). First we rewrite problem (21) with the notations (11) and (14):

$$
\begin{align*}
& \int_{\Omega}\left(b\left(x, \nabla u_{h}\right) \nabla u_{h} \cdot \nabla v_{h}+r\left(x, u_{h}\right) u_{h} v_{h}\right) d x+ \\
& +\int_{\Gamma} z\left(x, u_{h}\right) u_{h} v_{h} d \sigma=\int_{\Omega} \hat{f} v_{h} d x+\int_{\Gamma} \hat{\gamma} v_{h} d \sigma \tag{22}
\end{align*}
$$

$\forall v_{h} \in V_{h}^{0}$. We set

$$
\begin{equation*}
u_{h}=\sum_{j=1}^{\bar{n}} c_{j} \phi_{j} \tag{23}
\end{equation*}
$$

and look for the coefficients $c_{1}, \ldots, c_{\bar{n}}$. For any $\overline{\mathbf{c}}=\left(c_{1}, \ldots, c_{\bar{n}}\right) \in \mathbf{R}^{\bar{n}}, i=$ $1, \ldots, n$ and $j=1, \ldots, \bar{n}$, we set

$$
\begin{gather*}
b_{i j}(\overline{\mathbf{c}}):=\int_{\Omega} b\left(x, \sum_{k=1}^{\bar{n}} c_{k} \nabla \phi_{k}\right) \nabla \phi_{j} \cdot \nabla \phi_{i} d x \\
r_{i j}(\overline{\mathbf{c}}):=\int_{\Omega} r\left(x, \sum_{k=1}^{\bar{n}} c_{k} \phi_{k}\right) \phi_{j} \phi_{i} d x \\
z_{i j}(\overline{\mathbf{c}}):=\int_{\Gamma} z\left(x, \sum_{k=1}^{\bar{n}} c_{k} \phi_{k}\right) \phi_{j} \phi_{i} d \sigma, \quad d_{i}(\overline{\mathbf{c}}):=\int_{\Omega} \hat{f} \phi_{i} d x+\int_{\Gamma} \hat{\gamma} \phi_{i} d \sigma, \\
a_{i j}(\overline{\mathbf{c}}):=b_{i j}(\overline{\mathbf{c}})+r_{i j}(\overline{\mathbf{c}})+z_{i j}(\overline{\mathbf{c}}) . \tag{24}
\end{gather*}
$$

Putting (23) and $v_{h}=\phi_{i}$ into (22), we obtain the $n \times \bar{n}$ system of algebraic equations

$$
\begin{equation*}
\sum_{j=1}^{\bar{n}} a_{i j}(\overline{\mathbf{c}}) c_{j}=d_{i}, \quad i=1, \ldots, n \tag{25}
\end{equation*}
$$

Using the notations

$$
\begin{align*}
\mathbf{A}(\overline{\mathbf{c}}) & :=\left\{a_{i j}(\overline{\mathbf{c}})\right\}, \quad i, j=1, \ldots, n, \\
\tilde{\mathbf{A}}(\overline{\mathbf{c}}) & :=\left\{a_{i j}(\mathbf{c})\right\}, \quad i=1, \ldots, n ; j=n+1, \ldots, \bar{n},  \tag{26}\\
\mathbf{d} & :=\left\{d_{j}\right\}, \mathbf{c}:=\left\{c_{j}\right\}, \quad j=1, \ldots, n, \quad \text { and } \\
\tilde{\mathbf{c}} & :=\left\{c_{j}\right\}, \quad j=n+1, \ldots, \bar{n},
\end{align*}
$$

system (25) turns into

$$
\begin{equation*}
\mathbf{A}(\overline{\mathbf{c}}) \mathbf{c}+\tilde{\mathbf{A}}(\overline{\mathbf{c}}) \tilde{\mathbf{c}}=\mathbf{d} \tag{27}
\end{equation*}
$$

Defining further

$$
\overline{\mathbf{A}}(\overline{\mathbf{c}}):=\left[\begin{array}{ll}
\mathbf{A}(\overline{\mathbf{c}}) & \tilde{\mathbf{A}}(\overline{\mathbf{c}})
\end{array}\right], \quad \overline{\mathbf{c}}:=\left[\begin{array}{l}
\mathbf{c}  \tag{28}\\
\tilde{\mathbf{c}}
\end{array}\right],
$$

we rewrite (27) as follows

$$
\begin{equation*}
\overline{\mathbf{A}}(\overline{\mathbf{c}}) \overline{\mathbf{c}}=\mathbf{d} . \tag{29}
\end{equation*}
$$

In order to obtain a system with a square matrix, we enlarge our system to an $\bar{n} \times \bar{n}$ one. Since $u_{h}=g_{h}$ on $\partial \Omega$, the coordinates $c_{i}$ with $n+1 \leq i \leq \bar{n}$ satisfy automatically $c_{i}=g_{i}$, i.e.,

$$
\tilde{\mathbf{c}}=\tilde{\mathbf{g}}:=\left\{g_{j}\right\}, \quad j=n+1, \ldots, \bar{n},
$$

hence we can replace (27) by the equivalent system

$$
\left[\begin{array}{cc}
\mathbf{A}(\overline{\mathbf{c}}) & \tilde{\mathbf{A}}(\overline{\mathbf{c}})  \tag{30}\\
0 & \mathbf{I}
\end{array}\right]\left[\begin{array}{l}
\mathbf{c} \\
\tilde{\mathbf{c}}
\end{array}\right]=\left[\begin{array}{l}
\mathbf{d} \\
\tilde{\mathbf{g}}
\end{array}\right] .
$$

## 3 Maximum principle for the discretized problem

### 3.1 Background

First we recall a basic definition in the study of DMP (cf. [29, p. 23]):
Definition 2. A square $n \times n$ matrix $\mathbf{M}=\left(m_{i j}\right)_{i, j=1}^{n}$ is called irreducibly diagonally dominant if it satisfies the following conditions:
(i) $\mathbf{M}$ is irreducible, i.e., for any $i \neq j$ there exists a sequence of nonzero entries $\left\{m_{i, i_{1}}, m_{i_{1}, i_{2}}, \ldots, m_{i_{s}, j}\right\}$ of $M$, where $i, i_{1}, i_{2}, \ldots, i_{s}, j$ are distinct indices,
(ii) $\mathbf{M}$ is diagonally dominant, i.e., $\left|m_{i i}\right| \geq \sum_{\substack{j=1 \\ j \neq i}}^{n}\left|m_{i j}\right|, i=1, \ldots, n$,
(iii) for at least one index $i_{0} \in\{1, \ldots, n\}$ the above inequality is strict, i.e.,

$$
\left|m_{i_{0}, i_{0}}\right|>\sum_{\substack{j=1 \\ j \neq i_{0}}}^{n}\left|m_{i_{0}, j}\right| .
$$

Let us now consider a system of equations of order $(n+m) \times(n+m)$ :

$$
\overline{\mathbf{A}} \overline{\mathbf{c}}=\overline{\mathbf{b}},
$$

where the matrix $\overline{\mathbf{A}}$ has the following structure:

$$
\overline{\mathbf{A}}=\left[\begin{array}{cc}
\mathbf{A} & \tilde{\mathbf{A}}  \tag{31}\\
0 & \mathbf{I}
\end{array}\right] .
$$

Here $\mathbf{I}$ is the $m \times m$ identity matrix and $\mathbf{0}$ is the $m \times n$ zero matrix. Following [6], we introduce

Definition 3. An $(n+m) \times(n+m)$ matrix $\overline{\mathbf{A}}$ with the structure (31) is said to be of generalized nonnegative type if the following properties hold:
(i) $a_{i i}>0, \quad i=1, \ldots, n$,
(ii) $a_{i j} \leq 0, \quad i=1, \ldots, n, j=1, \ldots, n+m \quad(i \neq j)$,
(iii) $\sum_{j=1}^{n+m} a_{i j} \geq 0, \quad i=1, \ldots, n$,
(iv) There exists an index $i_{0} \in\{1, \ldots, n\}$ for which $\sum_{j=1}^{n} a_{i_{0}, j}>0$.
(v) $\mathbf{A}$ is irreducible.

Remark 3. In the original definition in [6, p. 343], it is assumed instead of the above properties (iv)-(v) that the principal block $\mathbf{A}$ is irreducibly diagonally dominant. However, the latter follows directly from Definition 3 under the given sign conditions on $a_{i j}$.

We also note that a well-known theorem [29, p. 85] implies in this case that $\mathbf{A}^{\mathbf{- 1}}>0$, i.e., the entries of the matrix $\mathbf{A}^{\mathbf{- 1}}$ are positive.

The known results on various discrete maximum principles (e.g., $[6,7,14$, 20]) are essentially based on the following theorem:

Theorem 5. Let $\overline{\mathbf{A}}$ be a $(n+m) \times(n+m)$ matrix with the structure (31), and assume that $\overline{\mathbf{A}}$ is of generalized nonnegative type in the sense of Definition 3.

If the vector $\overline{\mathbf{c}}=\left(c_{1}, \ldots, c_{n+m}\right) \in \mathbf{R}^{n+m}$ is such that $(\overline{\mathbf{A}} \overline{\mathbf{c}})_{i} \leq 0, i=1, \ldots, n$, then

$$
\begin{equation*}
\max _{i=1, \ldots, n+m} c_{i} \leq \max \left\{0, \max _{i=n+1, \ldots, n+m} c_{i}\right\} \tag{32}
\end{equation*}
$$

If, in addition,

$$
\begin{equation*}
\sum_{j=1}^{n+m} a_{i j}=0, \quad i=1, \ldots, n \tag{33}
\end{equation*}
$$

then

$$
\begin{equation*}
\max _{i=1, \ldots, n+m} c_{i}=\max _{i=n+1, \ldots, n+m} c_{i} \tag{34}
\end{equation*}
$$

Proof. As stated in Remark 3, A is irreducibly diagonally dominant. This, together with (i)-(iii), implies both statements (32) and (34), see [6, Th. 3] and [14, Th. 3], respectively.

### 3.2 Algebraic conditions for the discrete maximum principle

The following theorem is the main result of the present paper since it will allow us to derive various forms of the discrete maximum principle. The sign condition (36) is similar to the one given in [7, 14].

Theorem 6. Let Assumptions 2.1 hold and let

$$
\begin{equation*}
f(x)-q(x, 0) \leq 0, x \in \Omega, \quad \text { and } \quad \gamma(x)-s(x, 0) \leq 0, x \in \Gamma . \tag{35}
\end{equation*}
$$

Let us consider a family of simplicial triangulations $\mathcal{T}_{h}(h>0)$ satisfying the following property: for any $i=1, \ldots, n, j=1, \ldots, \bar{n}(i \neq j)$

$$
\begin{equation*}
\nabla \phi_{i} \cdot \nabla \phi_{j} \leq-\frac{\sigma_{0}}{h^{2}}<0 \tag{36}
\end{equation*}
$$

on supp $\phi_{i} \cap \operatorname{supp} \phi_{j}$ with $\sigma_{0}>0$ independent of $i, j$ and $h$.
(1) Let the triangulations $\mathcal{T}_{h}$ be regular, i.e., there exist constants $c_{1}, c_{2}>0$ such that for any $h>0$ and any simplex $T \in \mathcal{T}_{h}$

$$
\begin{equation*}
c_{1} h^{d} \leq \operatorname{meas}(T) \leq c_{2} h^{d} \tag{37}
\end{equation*}
$$

(where meas denotes d-dimensional measure). Then for sufficiently small $h$, the matrix $\overline{\mathbf{A}}(\overline{\mathbf{c}})$ defined in (28) is of generalized nonnegative type in the sense of Definition 3.
(2) More generally, for statement (1) to hold, it suffices to assume instead of (37) that the triangulations $\mathcal{T}_{h}$ are only quasi-regular in the following sense: the left-hand side of (37) is replaced by

$$
\begin{equation*}
c_{1} h^{\gamma} \leq \operatorname{meas}(T) \tag{38}
\end{equation*}
$$

with some $\gamma \geq d$ satisfying

$$
\begin{equation*}
2 \leq \gamma<3 \quad \text { if } d=2, \quad 3 \leq \gamma<\min \left\{\frac{12}{p_{1}-2}, 5-\frac{p_{2}}{2}\right\} \quad \text { if } d=3 \tag{39}
\end{equation*}
$$

(or in general, $d \leq \gamma<\min \left\{\frac{4 d}{\left(p_{1}-2\right)(d-2)}, 3+\frac{\left(4-p_{2}\right)(d-2)}{2}\right\}$ if $d \geq 3$ ) where $p_{1}$ and $p_{2}$ are defined in Assumptions 2.1, (A4).

Proof. The coefficients of $\overline{\mathbf{A}}(\overline{\mathbf{c}})$ satisfy

$$
a_{i j}(\overline{\mathbf{c}})=\int_{\Omega}\left[b\left(x, \nabla u_{h}\right) \nabla \phi_{i} \cdot \nabla \phi_{j}+r\left(x, u_{h}\right) \phi_{i} \phi_{j}\right] d x+\int_{\Gamma} z\left(x, u_{h}\right) \phi_{i} \phi_{j} d \sigma
$$

$(i=1, \ldots, n, j=1, \ldots, \bar{n})$. We now prove the properties (i)-(v) in the case (2); the conditions (38)-(39) are only used in part (ii).
(i) From our assumptions $b \geq \mu_{0}>0, r \geq 0$ and $z \geq 0$ we have

$$
a_{i i}(\overline{\mathbf{c}}) \geq \mu_{0} \int_{\Omega}\left|\nabla \phi_{i}\right|^{2} d x>0
$$

(ii) Let $i=1, \ldots, n, j=1, \ldots, \bar{n}$ with $i \neq j$ and let $\Omega_{i j}$ denote the interior of supp $\phi_{i} \cap \operatorname{supp} \phi_{j}$. If $\Omega_{i j}=\emptyset$ then

$$
a_{i j}(\overline{\mathbf{c}})=0 .
$$

If $\Omega_{i j} \neq \emptyset$ then properties (12) and (36) and the fact $0 \leq \phi_{i} \leq 1, i=1, \ldots, \bar{n}$, imply

$$
\begin{equation*}
a_{i j}(\overline{\mathbf{c}}) \leq-\frac{\sigma_{0}}{h^{2}} \mu_{0} \operatorname{meas}\left(\Omega_{i j}\right)+\int_{\Omega_{i j}} r\left(x, u_{h}\right) d x+\int_{\Gamma_{i j}} z\left(x, u_{h}\right) d \sigma, \tag{40}
\end{equation*}
$$

using notation $\Gamma_{i j}=\Gamma \cap \bar{\Omega}_{i j}$. Here, from (11) and Assumption (A4),

$$
\begin{aligned}
\int_{\Omega_{i j}} r\left(x, u_{h}\right) d x= & \int_{\Omega_{i j}} \frac{\partial q}{\partial \xi}\left(x, \theta u_{h}\right) d x \leq \int_{\Omega_{i j}}\left(\alpha_{1}(x)+\beta\left|\theta u_{h}\right|^{p_{1}-2}\right) d x \leq \\
& \leq \int_{\Omega_{i j}} \alpha_{1}(x) d x+\beta \int_{\Omega_{i j}}\left|u_{h}\right|^{p_{1}-2} d x
\end{aligned}
$$

(where we had some $\theta=\theta(x) \in[0,1]$ ), and in just the same way we have

$$
\int_{\Gamma_{i j}} z\left(x, u_{h}\right) d \sigma \leq \int_{\Gamma_{i j}} \alpha_{2}(x) d \sigma+\beta \int_{\Gamma_{i j}}\left|u_{h}\right|^{p_{2}-2} d \sigma
$$

Now we can estimate the integrals $\int_{\Omega_{i j}}\left|u_{h}\right|^{p_{1}-2} d x$ and $\int_{\Gamma_{i j}}\left|u_{h}\right|^{p_{2}-2} d \sigma$ as follows. We define $p^{*}:=\frac{2 d}{d-2}$ and $p^{* *}:=\frac{2(d-1)}{d-2}$ if $d \geq 3$, and $p^{*}:=p^{* *}:=+\infty$ if $d=2$. Then the Sobolev embedding estimates

$$
\begin{equation*}
\|v\|_{L^{p^{*}}(\Omega)} \leq k_{1}\|v\|_{1}, \quad\|v\|_{L^{p^{* *}}(\Gamma)} \leq k_{2}\|v\|_{1}, \quad v \in H^{1}(\Omega) \tag{41}
\end{equation*}
$$

hold with constants $k_{1}, k_{2}>0$, where $\|v\|_{1}=\|v\|_{H^{1}(\Omega)}$ (see [1]). Assume for a while that $p_{1}, p_{2}>2$ and let us fix real numbers $r$ and $t$ satisfying

$$
\begin{equation*}
\frac{\gamma}{2}<r \leq \frac{p^{*}}{p_{1}-2}, \quad \frac{d-1}{d+1-\gamma}<t \leq \frac{p^{* *}}{p_{2}-2} . \tag{42}
\end{equation*}
$$

Such numbers exist since for $d \geq 3$, by (39),

$$
\begin{aligned}
& \gamma<\frac{2 p^{*}}{p_{1}-2} \text { and } \quad \gamma<3+\frac{\left(4-p_{2}\right)(d-2)}{2}=d+1+\frac{\left(2-p_{2}\right)(d-2)}{2}= \\
&=d+1-\frac{\left(p_{2}-2\right)(d-1)}{p^{* *}}
\end{aligned}
$$

Further, $\gamma \geq 2$ implies $r \geq 1$ and $t \geq 1$. If $\frac{1}{r}+\frac{1}{s}=\frac{1}{t}+\frac{1}{l}=1$ then Hölder's inequality implies

$$
\begin{equation*}
\int_{\Omega_{i j}}\left|u_{h}\right|^{p_{1}-2} d x \leq\|1\|_{L^{s}\left(\Omega_{i j}\right)}\left\|\left|u_{h}\right|^{p_{1}-2}\right\|_{L^{r}\left(\Omega_{i j}\right)}=\operatorname{meas}\left(\Omega_{i j}\right)^{1 / s}\left\|u_{h}\right\|_{L^{\left(p_{1}-2\right) r}\left(\Omega_{i j}\right)}^{p_{1}-2} \tag{43}
\end{equation*}
$$

Here $\left(p_{1}-2\right) r \leq p^{*}$ and (41) imply

$$
\left\|u_{h}\right\|_{L^{\left(p_{1}-2\right) r}\left(\Omega_{i j}\right)}^{p_{1}-2} \leq\left\|u_{h}\right\|_{L^{\left(p_{1}-2\right) r}(\Omega)}^{p_{1}-2} \leq \text { const. } \cdot\left\|u_{h}\right\|_{L^{p^{*}}(\Omega)}^{p_{1}-2} \leq \text { const. } \cdot\left\|u_{h}\right\|_{1}^{p_{1}-2}
$$

Owing to the basic FEM convergence result, we have $\left\|u_{h}\right\|_{1} \rightarrow\left\|u^{*}\right\|_{1}$, where $u^{*}$ is the exact weak solution of our problem. Hence if $h$ is less than some fixed $h_{0}$ then (43) finally turns into

$$
\begin{equation*}
\int_{\Omega_{i j}}\left|u_{h}\right|^{p_{1}-2} d x \leq K_{1} \operatorname{meas}\left(\Omega_{i j}\right)^{1 / s} \tag{44}
\end{equation*}
$$

with some constant $K_{1}>0$ independent of $h$. In just the same way we obtain

$$
\begin{equation*}
\int_{\Gamma_{i j}}\left|u_{h}\right|^{p_{2}-2} d x \leq K_{2} \operatorname{meas}\left(\Gamma_{i j}\right)^{1 / l} \tag{45}
\end{equation*}
$$

Finally, if $p_{1}$ or $p_{2}$ equals 2 then the corresponding equality (44) or (45) holds with $s=1$ or $l=1$, respectively.

The integrals of $\alpha_{1}(x)$ and $\alpha_{2}(x)$ can be estimated with Hölder's inequality similarly to (43) by letting $\frac{2}{d}+\frac{1}{s^{\prime}}=\frac{1}{d-1}+\frac{1}{l^{\prime}}=1$ :

$$
\int_{\Omega_{i j}} \alpha_{1}(x) d x \leq K_{3} \operatorname{meas}\left(\Omega_{i j}\right)^{1 / s^{\prime}}, \quad \int_{\Gamma_{i j}} \alpha_{2}(x) d \sigma \leq K_{4} \operatorname{meas}\left(\Gamma_{i j}\right)^{1 / l^{\prime}}
$$

with $K_{3}=\left\|\alpha_{1}\right\|_{L^{d / 2}(\Omega)}$ and $K_{4}=\left\|\alpha_{2}\right\|_{L^{d-1}(\Gamma)}$.
Substituting all the estimates in (40), we obtain

$$
\begin{align*}
a_{i j}(\overline{\mathbf{c}}) \leq & -\frac{\sigma_{0} \mu_{0}}{h^{2}} \operatorname{meas}\left(\Omega_{i j}\right)+\beta K_{1} \operatorname{meas}\left(\Omega_{i j}\right)^{1 / s}+K_{3} \operatorname{meas}\left(\Omega_{i j}\right)^{1 / s^{\prime}}  \tag{46}\\
& +\beta K_{2} \operatorname{meas}\left(\Gamma_{i j}\right)^{1 / l}+K_{4} \operatorname{meas}\left(\Gamma_{i j}\right)^{1 / l^{\prime}}
\end{align*}
$$

We can write

$$
a_{i j}(\overline{\mathbf{c}}) \leq A_{1}^{i j}(h)+A_{2}^{i j}(h)+A_{3}^{i j}(h)+A_{4}^{i j}(h)
$$

where, with suitable constants $C_{0}, C_{1}, C_{2}, C_{3}, C_{4}>0$ independent of $h$ and $i, j$,

$$
\begin{aligned}
A_{1}^{i j}(h) & :=-\frac{C_{0}}{h^{2}} \operatorname{meas}\left(\Omega_{i j}\right)+C_{1} \operatorname{meas}\left(\Omega_{i j}\right)^{1 / s} \\
A_{2}^{i j}(h) & :=-\frac{C_{0}}{h^{2}} \operatorname{meas}\left(\Omega_{i j}\right)+C_{2} \operatorname{meas}\left(\Gamma_{i j}\right)^{1 / l} \\
A_{3}^{i j}(h) & :=-\frac{C_{0}}{h^{2}} \operatorname{meas}\left(\Omega_{i j}\right)+C_{3} \operatorname{meas}\left(\Omega_{i j}\right)^{1 / s^{\prime}}, \\
A_{4}^{i j}(h) & :=-\frac{C_{0}}{h^{2}} \operatorname{meas}\left(\Omega_{i j}\right)+C_{4} \operatorname{meas}\left(\Gamma_{i j}\right)^{1 / l^{\prime}} .
\end{aligned}
$$

We verify that for small enough $h$ we have $A_{k}^{i j}(h)<0(k=1,2,3,4)$.
Using $\frac{1}{r}+\frac{1}{s}=1$ and (38), we have

$$
\begin{aligned}
A_{1}^{i j}(h) & =\operatorname{meas}\left(\Omega_{i j}\right)^{1 / s}\left(-\frac{C_{0}}{h^{2}} \operatorname{meas}\left(\Omega_{i j}\right)^{1 / r}+C_{1}\right) \leq \\
& \leq \operatorname{meas}\left(\Omega_{i j}\right)^{1 / s}\left(-C_{5} h^{-2+(\gamma / r)}+C_{1}\right)
\end{aligned}
$$

Since (42) implies $\frac{\gamma}{r}<2$, the term in brackets tends to $-\infty$ as $h \rightarrow 0$ and hence $A_{1}^{i j}(h)<0$ for small $h$.

Using (38) again and the fact that meas $\left(\Gamma_{i j}\right) \leq$ const $\cdot h^{d-1} \quad$ (since $h$ is the diameter of the simplices and $\Gamma_{i j}$ lies on the $(d-1)$-dimensional surface), we have

$$
A_{2}^{i j}(h) \leq-C_{6} h^{\gamma-2}+C_{7} h^{\frac{d-1}{l}} .
$$

Since (42) implies $1-\frac{1}{l}=\frac{1}{t}<\frac{d+1-\gamma}{d-1}=1-\frac{\gamma-2}{d-1}$, we obtain $\frac{d-1}{l}>\gamma-2$, i.e., the second term tends to 0 faster and hence $A_{2}^{i j}(h)<0$ for small $h$.

The terms $A_{3}^{i j}(h)$ and $A_{4}^{i j}(h)$ can be handled similarly, since $s^{\prime}$ and $l^{\prime}$ satisfy the same estimates as $s$ and $l$. Namely, we have $\frac{d}{2}=\frac{p^{*}}{p^{*}-2}$ and $d-1=$ $\frac{p^{* *}}{p^{* *}-2}$, hence by substituting $\frac{d}{2}$ and $d-1$ for $r$ and $t$, respectively, we obtain that (42) holds in the special case $p_{1}=p^{*}$ and $p_{2}=p^{* *}$. Owing to the condition $\frac{2}{d}+\frac{1}{s^{\prime}}=\frac{1}{d-1}+\frac{1}{l^{\prime}}=1$, the numbers $s^{\prime}$ and $l^{\prime}$ play the same role as $s$ and $l$ and therefore the above estimates on $A_{1}^{i j}(h)$ and $A_{2}^{i j}(h)$ can be repeated for $A_{3}^{i j}(h)$ and $A_{4}^{i j}(h)$.

Altogether, we obtain that for small enough $h, A_{k}^{i j}(h)<0(k=1,2,3,4)$, that is, there exists $h_{0}>0$ such that

$$
\begin{equation*}
a_{i j}(\overline{\mathbf{c}})<0 \tag{47}
\end{equation*}
$$

for all $h \leq h_{0}$ and all $i, j$.
(iii) For any $i=1, \ldots, n$,

$$
\begin{align*}
\sum_{j=1}^{\bar{n}} a_{i j}(\overline{\mathbf{c}})= & \int_{\Omega}\left[b\left(x, \nabla u_{h}\right) \nabla \phi_{i} \cdot \nabla\left(\sum_{j=1}^{\bar{n}} \phi_{j}\right)+r\left(x, u_{h}\right) \phi_{i}\left(\sum_{j=1}^{\bar{n}} \phi_{j}\right)\right] d x \\
& +\int_{\Gamma} z\left(x, u_{h}\right) \phi_{i}\left(\sum_{j=1}^{\bar{n}} \phi_{j}\right) d \sigma  \tag{48}\\
= & \int_{\Omega} r\left(x, u_{h}\right) \phi_{i} d x+\int_{\Gamma} z\left(x, u_{h}\right) \phi_{i} d \sigma \quad \geq 0
\end{align*}
$$

using the fact that $\sum_{j=1}^{\bar{n}} \phi_{j} \equiv 1$ and $r, z, \phi_{i}$ are nonnegative.
(iv) Assume for contradiction that $\sum_{j=1}^{n} a_{i j}(\overline{\mathbf{c}})=0$ for all $i=1, \ldots, n$. This means that $\mathbf{A}(\overline{\mathbf{c}})$ carries the $n$-tuple of ones $\{1, \ldots, 1\}$ into the zero vector. This is impossible since $\mathbf{A}(\overline{\mathbf{c}})$ is symmetric and positive definite, and hence one-to-one.
(v) For any $i, j=1, \ldots, n$ with $i \neq j$, let us pick a sequence of neighbouring vertices $B_{i_{k}}(k=1, \ldots, s)$ in $\Omega$ that connect $B_{i}$ with $B_{j}$ (i.e., $i_{0}=i$ and $i_{s}=j$ ). Here (47) shows that $a_{i_{k}, i_{k+1}}(\overline{\mathbf{c}})<0$, hence by Definition 2, $\mathbf{A}(\overline{\mathbf{c}})$ is irreducible.

Theorem 6 enables us to derive the discrete the maximum principle for system (27):

Theorem 7. Under the conditions of Theorem 6, we have

$$
\begin{equation*}
\max _{\bar{\Omega}} u_{h} \leq \max \left\{0, \max _{\partial \Omega} g_{h}\right\} \tag{49}
\end{equation*}
$$

In particular, if $g \geq 0$, then $\max _{\bar{\Omega}} u_{h}=\max _{\partial \Omega} g_{h}$, and if $g \leq 0$, then we have the nonpositivity property $\max _{\bar{\Omega}} u_{h} \leq 0$.

Proof. Theorem 6 states that the condition of Theorem 5 is satisfied with $\overline{\mathbf{A}}(\mathbf{c})$ and $\bar{n}$ substituted for $\overline{\mathbf{A}}$ and $n+m$, respectively. Hence (32) yields

$$
\begin{equation*}
\max _{i=1, \ldots, \bar{n}} c_{i} \leq \max \left\{0, \max _{i=n+1, \ldots, \bar{n}} c_{i}\right\} \tag{50}
\end{equation*}
$$

Since $c_{i}=g_{i}$ for all $i=n+1, \ldots, \bar{n}$, estimate (50) is equivalent to (49).
The analogous minimum principle for system (27) can be verified in the same way.

Theorem 8. Let the conditions of Theorem 6 hold, except for (35) which is now replaced by

$$
\begin{equation*}
f(x)-q(x, 0) \geq 0 \quad(x \in \Omega) \quad \text { and } \quad \gamma(x)-s(x, 0) \geq 0 \quad(x \in \Gamma) \tag{51}
\end{equation*}
$$

Then we have

$$
\begin{equation*}
\min _{\bar{\Omega}} u_{h} \geq \min \left\{0, \min _{\partial \Omega} g_{h}\right\} \tag{52}
\end{equation*}
$$

In particular, if $g \leq 0$, then $\min _{\bar{\Omega}} u_{h}=\min _{\partial \Omega} g_{h}$, and if $g \geq 0$, then we have the nonnegativity property $\min _{\bar{\Omega}} u_{h} \geq 0$.

Let us now consider the special case $q \equiv 0$ and $s \equiv 0$. Then the counterpart of Theorem 3 is valid, which we now formulate for both the maximum and minimum principles. Moreover, the strict negativity in (36) can be replaced by the weaker nonnegativity property, regularity conditions on the mesh like (38)-(39) are not required, and the result for a proper mesh holds for all parameters $h$ instead of only sufficiently small $h$.

Theorem 9. Let us consider the following special case of problem (1):

$$
\left\{\begin{align*}
-\operatorname{div}(b(x, \nabla u) \nabla u) & =f(x) & & \text { in } \Omega \backslash \Gamma  \tag{53}\\
{[u]_{\Gamma} } & =0 & & \text { on } \Gamma \\
{\left[b(x, \nabla u) \frac{\partial u}{\partial \nu}\right]_{\Gamma} } & =\gamma(x) & & \text { on } \Gamma \\
u & =g(x) & & \text { on } \partial \Omega
\end{align*}\right.
$$

Let (A1)-(A3) of Assumptions 2.1 hold and let the triangulation $\mathcal{T}_{h}$ satisfy the following property: for any $i=1, \ldots, n, j=1, \ldots, \bar{n}(i \neq j)$

$$
\begin{equation*}
\nabla \phi_{i} \cdot \nabla \phi_{j} \leq 0 \tag{54}
\end{equation*}
$$

Then the following results hold:
(1) If $f \leq 0$ and $\gamma \leq 0$, then $\max _{\bar{\Omega}} u_{h}=\max _{\partial \Omega} g_{h}$.
(2) If $f \geq 0$ and $\gamma \geq 0$, then $\min _{\bar{\Omega}} u_{h}=\min _{\partial \Omega} g_{h}$.
(3) If $f=0$ and $\gamma=0$, then the ranges of $u_{h}$ and $g_{h}$ coincide, i.e., we have $\left[\min _{\bar{\Omega}} u_{h}, \max _{\bar{\Omega}} u_{h}\right]=\left[\min _{\partial \Omega} g_{h}, \max _{\partial \Omega} g_{h}\right]$ for the corresponding intervals.

Proof. (1) The conditions of Theorem 5 follow similarly as in Theorem 6. The difference arises in proving property (ii), i.e., $a_{i j}(\mathbf{c}) \leq 0$, where only (54) is sufficient, since the assumptions $q \equiv 0$ and $s \equiv 0$ imply $r \equiv 0$ and $z \equiv 0$. In order to apply statement (34) of Theorem 5 , it remains to verify that $\sum_{j=1}^{\bar{n}} a_{i j}(\mathbf{c})=0, i=1, \ldots, n$. Since $r \equiv 0$ and $z \equiv 0$, the argument used in (48) yields that this holds indeed. Statement (2) follows from (1) by replacing $u$ by $-u$, and (3) is a direct consequence of (1) and (2).

Remark 4. Conditions (36) and (54) can be in fact relaxed such that $\nabla \phi_{i}$. $\nabla \phi_{j}$ need not be negative resp. nonpositive on each element, see [14, Remark $6]$ for details.

### 3.3 Geometric conditions on the mesh

The conditions in the preceding subsection that guarantee the DMP have apparent geometric interpretations for our simplicial meshes. This relies on the fact that the values $\nabla \phi_{i} \cdot \nabla \phi_{j}$ are constant on each simplicial element, hence conditions (36) and (54) are not very difficult to check. Indeed, it is shown in [3] that

$$
\begin{equation*}
\left.\nabla \phi_{i} \cdot \nabla \phi_{j}\right|_{T}=-\frac{\operatorname{meas}_{d-1}\left(S_{i}\right) \cdot \operatorname{meas}_{d-1}\left(S_{j}\right)}{d^{2}\left(\operatorname{meas}_{d}(T)\right)^{2}} \cos \left(S_{i}, S_{j}\right) \quad \text { for } i \neq j \tag{55}
\end{equation*}
$$

where $T$ is a $d$-dimensional simplex with vertices $P_{1}, \ldots, P_{d+1}, S_{i}$ is the face of $T$ opposite to $P_{i}$, and $\cos \left(S_{i}, S_{j}\right)$ is the cosine of the interior angle between faces $S_{i}$ and $S_{j}$.

Thus, in order to satisfy condition (36) or (54), it is sufficient if the employed simplicial mesh is acute or nonobtuse, respectively (see [17, 18, 22], where also mesh refinement procedures preserving the above-mentioned geometrical properties are presented). We note that the conditions of acuteness or nonobtuseness are sufficient but not necessary: as referred to in Remark 4, the DMP may still hold if some obtuse interior angles occur in the simplices of the meshes. This is analogous to the case of linear problems [19, 28].

We note that the results can be easily extended to the case of meshes consisting of block elements, treated as in [15, Sect. 5.2].

### 3.4 Some applications to model problems

We quote two examples of problems where suitable discrete nonnegativity or nonpositivity properties are valid.

### 3.4.1 Semilinear equations: reaction-diffusion problems with localized autocatalytic chemical reactions

The problem

$$
\left\{\begin{align*}
-\Delta u & =f(x) & & \text { in } \Omega \backslash \Gamma,  \tag{56}\\
{[u]_{\Gamma} } & =0 & & \text { on } \Gamma, \\
{\left[\frac{\partial u}{\partial \nu}\right]_{\Gamma}+s(x, u) } & =0 & & \text { on } \Gamma, \\
u & =0 & & \text { on } \partial \Omega,
\end{align*}\right.
$$

in a planar domain $\Omega \subset \mathbf{R}^{2}$ describes a chemical reaction-diffusion process where the reaction is localized at the curve $\Gamma$, further, the reaction is autocatalytic, i.e., the growth of the concentration $u \geq 0$ speeds up the rate of the reaction, that is $\frac{\partial s(x, u)}{\partial u} \geq 0$ (see, e.g., $[12,13]$ ). The reaction function $s$ grows at most polynomially in $u$, hence Assumptions 2.1 hold. The fact that there is no reaction without material is expressed by $s(x, 0)=0$, further, we may assume that the source term $f$ is nonnegative. These conditions imply that the requirement $u \geq 0$ is satisfied, see subsection 2.3, moreover, the boundary conditions yield $\min _{\bar{\Omega}} u=0$. As a special case of Theorem 8 , we obtain the corresponding discrete minimum principle:

Corollary 1. Let $u_{h}$ be the FEM solution to problem (56) under a FEM discretization with the acuteness property (36). If $h$ is sufficiently small then

$$
\min _{\bar{\Omega}} u_{h}=0 .
$$

### 3.4.2 Linear equations

The following linear interface model arises in many applications such as biochemistry or multiphase flow, see, e.g., [26]:

$$
\left\{\begin{align*}
-\operatorname{div}(k(x) \nabla u) & =f(x) & & \text { in } \Omega \backslash \Gamma,  \tag{57}\\
{[u]_{\Gamma} } & =0 & & \text { on } \Gamma, \\
{\left[k(x) \frac{\partial u}{\partial \nu}\right]_{\Gamma} } & =\gamma(x) & & \text { on } \Gamma, \\
u & =0 & & \text { on } \partial \Omega,
\end{align*}\right.
$$

where the bounded measurable function $k$ is discontinuous on $\Gamma$. In addition, it suffices to assume that $k$ has a positive lower bound and $f \in L^{2}(\Omega)$, $\gamma \in L^{2}(\Gamma)$. Then, as a special case of Theorem 9 , we obtain the corresponding discrete maximum and minimum principles:

Corollary 2. Let $u_{h}$ be the FEM solution to problem (57) under a FEM discretization with the nonobtuseness property (54).

If $f \leq 0$ and $\gamma \leq 0$ then $\max _{\bar{\Omega}} u_{h}=0$, and if $f \geq 0$ and $\gamma \geq 0$ then $\min _{\bar{\Omega}} u_{h}=0$.

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(continued from the back cover)
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